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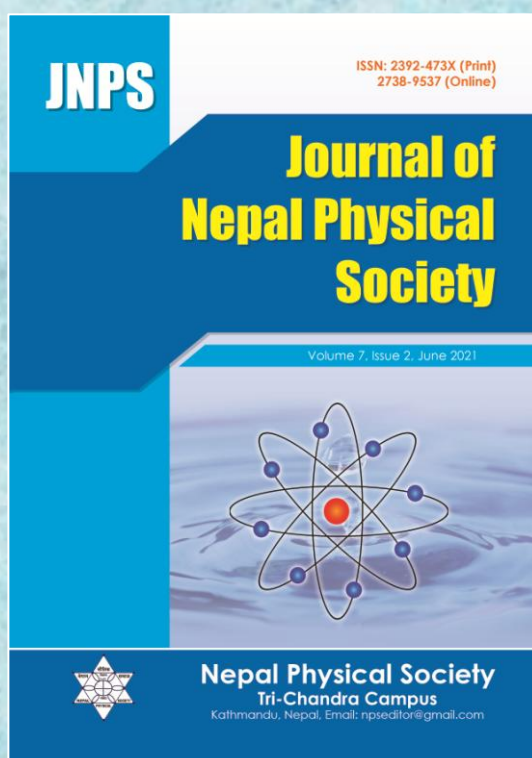
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# Structural, Electronic and Magnetic Properties of Impurities Defected Graphene/MoS<sub>2</sub> Van Der Waals Heterostructure: First-principles Study

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## ABSTRACT

Two-dimensional (2D) pristine and defected van der Waals (vdW) heterostructure (HS) materials open up fortune in nanoelectronic and optoelectronic devices. So, they are compatible for designing in the fields of device applications. In the present work, we studied structural, electronic and magnetic properties of vdW (HS) graphene/MoS<sub>2</sub> ((HS)G/MoS<sub>2</sub>), Nb impurity defect in vdW (HS) graphene/MoS<sub>2</sub> (Nb-(HS)G/MoS<sub>2</sub>), and Tc impurity defect in vdW (HS) graphene/MoS<sub>2</sub> (Tc-(HS)G/MoS<sub>2</sub>) materials by using spin-polarized DFT-D2 method. We examined the structure of these materials, and found that they are stable. Based on band structure analysis, we found that (HS)G/MoS<sub>2</sub>, Nb-(HS)G/MoS<sub>2</sub> and Tc-(HS)G/MoS<sub>2</sub> have metallic characteristics. Also, (HS)G/MoS<sub>2</sub> and Tc-(HS)G/MoS<sub>2</sub> materials have n-type Schottky contact, while Nb-(HS)G/MoS<sub>2</sub> material has p-type Schottky contact. To understand the magnetic properties of materials, we have used DoS, IDoS and PDoS calculations. We found that (HS)G/MoS<sub>2</sub> is a non-magnetic material, but Nb-(HS)G/MoS<sub>2</sub> and Tc-(HS)G/MoS<sub>2</sub> are magnetic materials. Magnetic moment of Nb-(HS)G/MoS<sub>2</sub> and Tc-(HS)G/MoS<sub>2</sub> materials are  $-0.24 \mu_B/\text{cell}$  and  $+0.07 \mu_B/\text{cell}$  values respectively from DoS/PDoS calculations, and  $0.26 \mu_B/\text{cell}$  and  $0.08 \mu_B/\text{cell}$  values respectively from IDoS calculations. Up-spin and down-spin states of electrons in 2p orbital of C atoms, 3p orbital of S atoms, 4d orbital of Mo atoms, 4d orbital of Tc atom in Tc-(HS)G/MoS<sub>2</sub>, and 2p orbital of C atoms, 3p orbital of S atoms, 4p & 4d orbitals of Mo atoms, 4p & 4d orbitals of Nb atom in Nb-(HS)G/MoS<sub>2</sub> have major contribution for the development of magnetic moment.

**Keywords:** DFT-D2; Heterostructure; Impurity defects; Magnetic moment; Spin states.

## 1. INTRODUCTION

Graphene is two dimensional (2D) stretchable,  $sp^2$ -hybridized, zero band gap and honeycomb lattice structure of carbon atoms. It has good electronic properties such as, high carrier mobility and high thermal conductivity, observable Quantum Hall effect at room temperature and existence of two dimensional gases of massless Dirac fermions [1-4]. Due to these properties, it has potential applications in the areas of nanoelectronic, optoelectronic, sensing and hydrogen storage devices [5-7]. Thus, graphene is a favorable material for the researchers to predict additional

desirable properties [4, 8]. Molybdenum disulphide (MoS<sub>2</sub>) is 2D transition metal dichalcogenide (TMD) wide band gap semiconductor of band gap 1.80 eV [9]. It can be constructed by assembling Mo and S atoms (i.e. S-Mo-S) in a triangular prismatic configuration [10]. MoS<sub>2</sub> has great applications in the fields of nanoelectronic and optoelectronic devices [11-14]. However, 2D materials, graphene due to zero bandgap energy, and monolayer MoS<sub>2</sub> due to wide gap energy have limited properties. Therefore, to tune the desirable properties in 2D materials, graphene is combined with other 2D materials and formed heterostructure

(HS). The graphene based 2D (HS) materials (G/MoS<sub>2</sub>, G/h-BN) are used to produce novel ways for engineering electronic and optoelectronic devices [15-17]. The (HS)G/MoS<sub>2</sub> (heterostructure of graphene with MoS<sub>2</sub>) offers mechanical, electronic, optical, and transport properties. Thus, it can be used in the fields of device applications [18]. Moreover, the substitution of any atom by foreign atom (impurity defect) in (HS)G/MoS<sub>2</sub> or removal of any atom (vacancy defect) from (HS)G/MoS<sub>2</sub> structure is one of the promising approaches to modify and exploit unwanted properties of any constituent. The impurity of Technetium (Tc) atom or Niobium (Nb) atom in (HS)G/MoS<sub>2</sub> changes its electronic and magnetic properties. Defects in (HS) give peculiar properties and carry out the new materials [19-22]. Hence, defects impact the properties of (HS) materials [23]. To our knowledge, effect of impurity defects on structural, electronic and magnetic properties of vdW (HS)G/MoS<sub>2</sub> material have not been reported yet. Therefore, in this paper, we learned the structural, electronic and magnetic properties of vdW (HS)G/MoS<sub>2</sub>, and effect of Nb impurity defect in vdW (HS)G/MoS<sub>2</sub> (Nb-(HS)G/MoS<sub>2</sub>) material & Tc impurity defects in vdW (HS)G/MoS<sub>2</sub> (Tc-(HS)G/MoS<sub>2</sub>) material by spin-polarized DFT-D2 method based on first-principles calculations. Figs. 1(a-c) illustrated pristine (HS)G/MoS<sub>2</sub>, Nb-(HS)G/MoS<sub>2</sub> and Tc-(HS)G/MoS<sub>2</sub> materials respectively.

## 2. METHODS AND MATERIALS

We used spin-polarized DFT-D2 method to learn the structural, electronic and magnetic properties of (HS)G/MoS<sub>2</sub> and impurity defects (Nb and Tc impurity defects) in vdW (HS)G/MoS<sub>2</sub> using Quantum ESPRESSO (QE) computational tool [24-26]. The electronic exchange and correlation effects in the systems are treated by Generalized Gradient Approximation (GGA) using Perdew-Burke-Ernzerhof (PBE) [27]. Rappe-Rabe-Kaxiraas-Joannopoulos (RRKJ) model of ultra-soft pseudo-potentials is used to deal the chemically active valence electrons in the systems. All structures are optimized and relaxed by Broyden-Fletcher-Goldfarb-Shanno (BFGS) [28] method. At first we have prepared G/MoS<sub>2</sub> heterostructure material by (4×4) supercell structure of graphene and (3×3) supercell structure of monolayer MoS<sub>2</sub> with 4.11% lattice mismatch. To avoid the physical interactions in the stacking direction of (HS)G/MoS<sub>2</sub> material, we kept a vacuum distance greater than 18 Å. The

Nb impurity defect in (HS)G/MoS<sub>2</sub> (i.e. Nb-(HS)G/MoS<sub>2</sub>) and Tc Impurity defect in (HS)G/MoS<sub>2</sub> (i.e. Tc-(HS)G/MoS<sub>2</sub>) are obtained by replacing centre Mo atom of (HS)G/MoS<sub>2</sub> by Nb atom and Tc atom respectively. Then, we obtained relax Nb-(HS)G/MoS<sub>2</sub> and Tc-(HS)G/MoS<sub>2</sub> structures by using BFGS method as shown in figs. 1(b-c). We performed self consistent total energy calculation after relax calculations. For this, the Brillouin-zone of (HS) is figured out in k-space using Monkhorst-Pack scheme [29] with suitable number of mesh (6×6×1) of k-points, which is determined from the convergence test. The plane-wave expansion with kinetic energy cut-off 35 Ry and charge density cut-off 350 Ry are used for wave function and charge density respectively. Kinetic energy cut-off and charge density cut-off values are calculated from convergence plot of total energy versus energy cut-off wave function. We used Marzari-Vanderbilt [30] method of smearing for occupations and 0.001 Ry value of degauss. We have taken “david” diagonalization method with “plain” mixing mode for self consistency with 0.6 mixing factor. Also, we have taken a mesh of (6×6×1) k-points for band structure calculations (100 k-points are chosen along the high symmetric points), and meshes of (12×12×1) k-points are used for density of states (DoS) and projected density of states (PDOS) calculations.

## 3. RESULTS AND DISCUSSION

In this section, we discussed in details about structural, electronic and magnetic properties of pristine (HS)G/MoS<sub>2</sub>, and Nb & Tc impurity defected (HS)G/MoS<sub>2</sub> materials based on spin-polarized DFT-D2 method.

### 3.1 Structural Analysis

We have prepared vertical stacking configuration of vdW (HS)G/MoS<sub>2</sub> material by keeping supercell of graphene and supercell of MoS<sub>2</sub>. The lattice mismatch is found to be 4.11% in vdW (HS)G/MoS<sub>2</sub>. Lattice mismatch is fixed by differing the lattice constant because there is no direct chemical bonding between the constituent layers. Stability of (HS)G/MoS<sub>2</sub> material is checked by its calculated binding energy. Binding energy (-0.40 eV) of it is obtained by using equation (1), [31];

$$E_b = E_{(\text{HS})\text{G}/\text{MoS}_2} - E_G - E_{\text{MoS}_2} \dots (1)$$

Where,  $E_{(\text{HS})\text{G}/\text{MoS}_2}$ ,  $E_G$  and  $E_{\text{MoS}_2}$  represent total ground state energy of (HS)G/MoS<sub>2</sub>, supercell of

graphene and supercell of MoS<sub>2</sub> respectively. The obtained binding energy (-0.40 eV) of (HS)G/MoS<sub>2</sub> indicates that material is energetically stable at ground state because negative value of binding energy of materials implies that they are energetically stable at ground state. We have also calculated the inter-layer distance the constituents of (HS) and found 3.37 Å value. Our calculated values of binding energy and inter-layer distance are comparable with other 2D VdW heterostructure materials [32]. In addition, we have created the stable Nb-(HS)G/MoS<sub>2</sub> and Tc-(HS)G/MoS<sub>2</sub> impurity defects materials as shown in figs. 1(b-c). The binding energy of Nb-(HS)G/MoS<sub>2</sub> is -0.35 eV and Tc-(HS)G/MoS<sub>2</sub> is -0.38 eV are calculated by using equation (2), [31]; because binding energy determines the stability of materials.

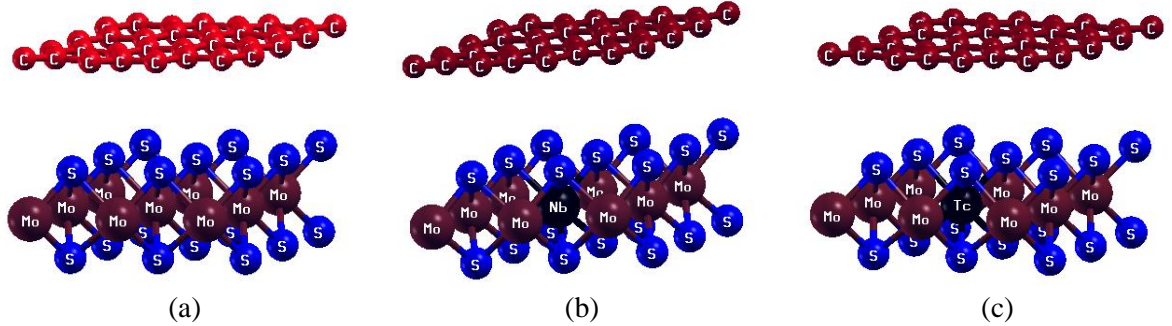
$$E_b = E_{(\text{HS-G/MoS}_2)_{i-d}} - E_G - E_{(\text{MoS}_2)_{i-d}} \dots (2)$$

Where,  $E_{(\text{HS-G/MoS}_2)_{i-d}}$ ,  $E_{(\text{MoS}_2)_{i-d}}$  &  $E_G$  represent total ground state energy of impurity (Nb atom or Tc atom) defected (HS)G/MoS<sub>2</sub>, impurity defected monolayer MoS<sub>2</sub> & graphene respectively. We measured the inter-layer distances between graphene and MoS<sub>2</sub> in Nb-(HS)G/MoS<sub>2</sub> and Tc-(HS)G/MoS<sub>2</sub> materials, and found that they are 3.54 Å and 3.42 Å respectively, which are comparable

to other 2D hetrostructure materials of values 3.31 Å and 3.19 Å [32]. Thus, the estimated binding energies and inter-layer distances of impurity defected materials are also comparable with other vdW heterostructure materials [32]. We can also test the stability of impurity defected materials relative to pristine material; we can calculate the defects formation energy ( $E_f$ ). Defects formation energy is obtained by using the values of total energy of impurity defects in (HS)G/MoS<sub>2</sub> ( $E_{t1}$ ), total energy of pristine (HS)G/MoS<sub>2</sub> ( $E_{t2}$ ), energy of impurity (Nb or Tc) atom ( $E_1$ ), and energy of isolated (replaced) Mo atom ( $E_2$ ) by equation (3), [33];

$$E_f = (E_{t1} - E_{t2}) - (E_1 - E_2) \dots (3)$$

Defects formation energy of Nb-(HS)G/MoS<sub>2</sub> and Tc-(HS)G/MoS<sub>2</sub> materials are found to be 2.54 eV and 2.56 eV respectively. The negative value of formation energy indicates that the system is more stable than the pristine one, while positive value of formation energy implies that the system requires external energy for the formation of impurity defects in material. Also, lower value of defects formation energy means, materials can be favorable in computation work.



**Fig. 1:** (a) Structure of pristine (HS)G/MoS<sub>2</sub> material, (b) Structure of Nb-(HS)G/MoS<sub>2</sub> material, (c) Structure of Tc-(HS)G/MoS<sub>2</sub> material.

From the structural analysis of pristine (HS)G/MoS<sub>2</sub> and impurity defects Nb-(HS)G/MoS<sub>2</sub> & Tc-(HS)G/MoS<sub>2</sub> materials, we found that all are stable and compactness of pristine has greater than that of impurity defects materials.

### 3.2 Electronic and Magnetic Properties

Solid state matter is rigid and its degree of rigidity depends upon the relative compactness of the constituent atoms. In an isolated atom, the electrons

occupy atomic orbitals each of which has a discrete energy level while atomic orbitals overlap when two identical atoms join together to form a molecule. In solid with N identical atoms, each atomic orbital splits into N discrete molecular orbitals, each with a different energy. In a macroscopic piece of solid, the number of atoms is large ( $N \approx 10^{23}$ ), so is the number of orbitals and is very closely spaced in energy ( $10^{-22}$  eV). Thus, they

form continuum of energy as ‘energy band’. Electronic properties of material can be predicted by the analysis of its band structure. In this section, we discuss electronic and magnetic properties of (HS)G/MoS<sub>2</sub>, Nb-(HS)G/MoS<sub>2</sub> and Tc-(HS)G/MoS<sub>2</sub> materials on the basis of their band structures and density of states (DoS), integrated density of states (IDoS) and projected density of states (PDoS) analysis. The band structures of (HS)G/MoS<sub>2</sub>, Nb-(HS)G/MoS<sub>2</sub> and Tc-(HS)G/MoS<sub>2</sub> materials are shown in fig. 2(a), fig. 3(a) and fig. 4(a) respectively. In band structures, x-axis represents  $\Gamma$ -centre symmetric points ( $\Gamma$ -M-K- $\Gamma$ ) in the first Brillouin-zone, where 100 k-points are kept along the particular direction of irreducible Brillouin-zone, and y-axis represents the corresponding energy values, also horizontal dotted line represents the Fermi energy level. In band structure of pristine (HS)G/MoS<sub>2</sub>, Dirac cone is formed in the conduction band at 0.56 eV distance from Fermi energy level. It means electrons flow from valence band to conduction band. Hence, (HS)G/MoS<sub>2</sub> has metallic properties. Also, n-type Schottky contact is formed with Schottky barrier height of 0.42 eV in (HS)G/MoS<sub>2</sub>. This value agrees with the reported value 0.49 eV [34, 35]. In band structures of impurity defects Nb-(HS)G/MoS<sub>2</sub> and Tc-(HS)G/MoS<sub>2</sub> materials, Dirac cone is formed in conduction band at 0.70 eV and 0.01 eV distances from the Fermi energy level as shown in fig. 3(a) and fig. 4(a) respectively. In both

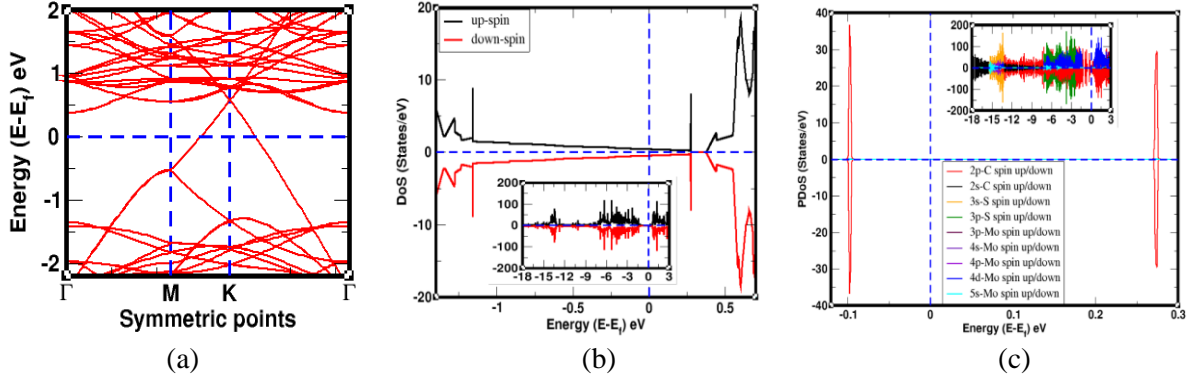
materials, electrons flow from valence band to conduction band, hence they have metallic properties. The n-type Schottky contact is still preserved in Tc-(HS)G/MoS<sub>2</sub>, while in Nb-(HS)G/MoS<sub>2</sub> material p-type Schottky contact is formed. Therefore, n-type Schottky contact of (HS)G/MoS<sub>2</sub> shifts to p-type Schottky contact due to Nb impurity atom in it. We know that electronic configurations of valence electrons in Mo, Tc, Nb, S & C atoms present in the materials are [Kr] 4d<sup>5</sup> 5s<sup>1</sup>, [Kr] 4d<sup>5</sup> 5s<sup>2</sup>, [Kr] 4d<sup>4</sup> 5s<sup>1</sup>, [Ne] 3s<sup>2</sup> 3p<sup>4</sup> & [He] 2s<sup>2</sup> 2p<sup>2</sup> respectively. Each Mo atom has one unpaired up-spin in 5s orbital and 4d<sub>xy</sub>, 4d<sub>xz</sub>, 4d<sub>yz</sub>, 4d<sub>x<sup>2</sup>-y<sup>2</sup></sub>, 4d<sub>z<sup>2</sup></sub> sub-orbitals; Tc atom has paired spins in 5s orbital and one unpaired up-spin in 4d<sub>xy</sub>, 4d<sub>xz</sub>, 4d<sub>yz</sub>, 4d<sub>x<sup>2</sup>-y<sup>2</sup></sub>, 4d<sub>z<sup>2</sup></sub> sub-orbitals; Nb atom has one unpaired up-spin in 5s orbital, 4d<sub>xy</sub>, 4d<sub>xz</sub>, 4d<sub>yz</sub>, 4d<sub>x<sup>2</sup>-y<sup>2</sup></sub> sub-orbitals and vacant in 4d<sub>z<sup>2</sup></sub> sub-orbital; S atom has paired spins (up-spin & down-spin) in 3p<sub>x</sub> sub-orbital and one unpaired up-spin in 3p<sub>y</sub>, 3p<sub>z</sub> sub-orbitals; C atom has single up-spin in 2p<sub>x</sub>, 2p<sub>y</sub> and vacant in 2p<sub>z</sub> sub-orbitals. By the configurations of unpaired (up-spin and down-spin) spins of electrons in the orbital of atoms in (HS)G/MoS<sub>2</sub>, Nb-(HS)G/MoS<sub>2</sub> and Tc-(HS)G/MoS<sub>2</sub> materials created different values of Fermi energy and Dirac (cone) point. The different values of Fermi energy, Dirac point, Dirac shift, Fermi shift, total energy, binding energy and inter-layer distances of above mentioned materials are given in table 1.

**Table 1: Fermi energy ( $E_f$ ), Fermi energy shift ( $E_s$ ), Dirac (point) cone ( $D_p$ ), Dirac (point) cone shift ( $D_s$ ), total energy ( $E_t$ ), binding energy ( $E_b$ ) and inter-layer distances ( $D_{l-d}$ ) of (HS)G/MoS<sub>2</sub>, Nb-(HS)G/MoS<sub>2</sub> and Tc-(HS)G/MoS<sub>2</sub> materials.**

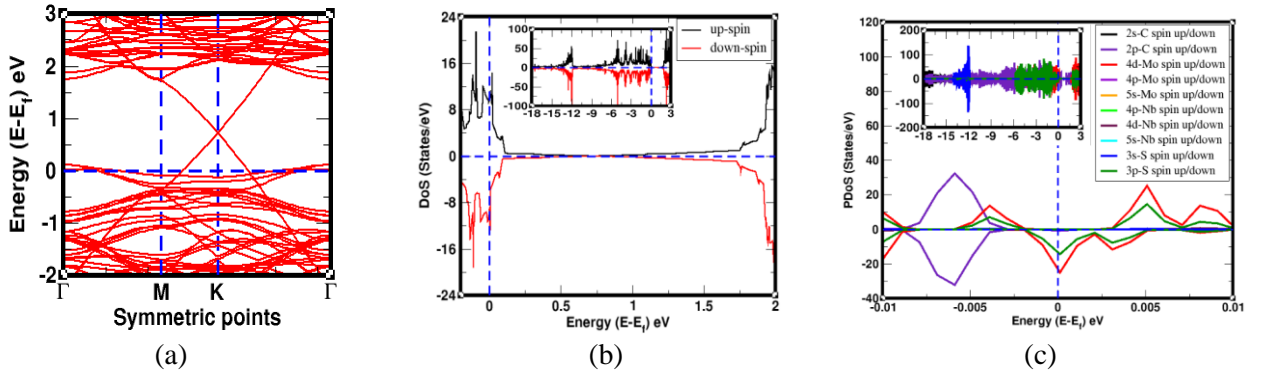
(HS)G/MoS <sub>2</sub> , Nb-(HS)G/MoS <sub>2</sub> and Tc-(HS)G/MoS <sub>2</sub> materials	$E_f$ (eV)	$E_s$ (eV)	$D_p$ (eV)	$D_s$ (eV)	$E_t$ (Ry)	$E_b$ (eV)	$D_{l-d}$ (Å)
(HS)G/MoS <sub>2</sub>	0.32	-	0.56	-	-2106.17	-0.40	3.37
Nb-(HS)G/MoS <sub>2</sub>	-0.50	-0.82	0.70	+0.14	-2076.54	-0.35	3.54
Tc-(HS)G/MoS <sub>2</sub>	0.78	+0.46	0.01	-0.55	-2138.61	-0.38	3.42

Based on DoS and PDoS calculations, we can investigate the magnetic properties of materials. Symmetrically distributed total up-spin and total down-spin in DoS/PDoS of materials means, materials has non-magnetic properties, and asymmetrically distributed unpaired up-spin and down-spin of electrons in the orbital of atom give magnetic properties of materials. The DoS and PDoS plots of (HS)G/MoS<sub>2</sub>, Nb-(HS)G/MoS<sub>2</sub> and

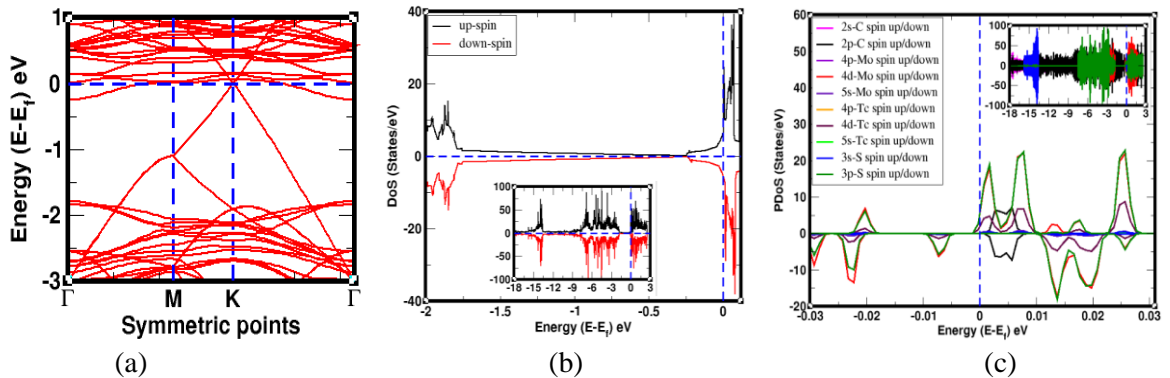
Tc-(HS)G/MoS<sub>2</sub> materials are given in figs 2(b-c), figs. 3(b-c) and figs. 4(b-c) respectively, where vertical dotted line distinguished energy bands (the region below zero scale is called valence band and the region above zero scale is called conduction band), and horizontal dotted line separates distributed up-spin and down-spin states of electrons in the orbitals of atoms in materials.



**Fig. 2:** (a) Band structure of (HS)G/MoS<sub>2</sub> material, (b) DoS of up-spin and down-spin states of electrons in the orbitals of C, Mo & S atoms of (HS)G/MoS<sub>2</sub> material (c) PDoS of individual up-spin and down-spin states of electrons in the orbitals of C, Mo & S atoms of (HS)G/MoS<sub>2</sub> material, where in band structure, horizontal dotted line represents Fermi level and in DoS/PDoS plots, vertical dotted line represents Fermi level. In all DoS and PDoS plots, insets represent the symmetrically distributed total up-spin and total down-spin of electrons in the orbital of atoms present in (HS)G/MoS<sub>2</sub> material.



**Fig. 3:** (a) Band structure of Nb-(HS)G/MoS<sub>2</sub> material, (b) DoS of up-spin and down-spin states of electrons in the orbitals of C, Mo, S & Nb atoms of Nb-(HS)G/MoS<sub>2</sub> material (c) PDoS of individual up-spin and down-spin states of electrons in the orbitals of C, Mo, S & Nb atoms of Nb-(HS)G/MoS<sub>2</sub> material, where in band structure, horizontal dotted line represents Fermi level and in DoS/PDoS plots, vertical dotted line represents Fermi level. In all DoS and PDoS plots, insets represent the asymmetrically distributed total up-spin and total down-spin of electrons in the orbital of atoms present in Nb-(HS)G/MoS<sub>2</sub> material.



**Fig. 4:** (a) Band structure of Tc-(HS)G/MoS<sub>2</sub> material, (b) DoS of up-spin and down-spin states of electrons in the orbitals of C, Mo, S & Tc atoms of Tc-(HS)G/MoS<sub>2</sub> material (c) PDoS of individual up-spin and down-spin states of electrons in the orbitals of C, Mo, S & Tc atoms of Tc-(HS)G/MoS<sub>2</sub> material, where in band structure, horizontal dotted line represents Fermi level and in DoS/PDoS plots, vertical dotted line represents Fermi level. In all DoS and PDoS plots, insets represent the asymmetrically distributed total up-spin and total down-spin of electrons in the orbital of atoms present in Tc-(HS)G/MoS<sub>2</sub> material.

The detail calculations of magnetic moment due to spin states of electrons in the orbitals of C, Mo, S, Nb & Tc

atoms in PDoS of (HS)G/MoS<sub>2</sub>, Nb-(HS)G/MoS<sub>2</sub> and Tc-(HS)G/MoS<sub>2</sub> materials are given in table 2.

**Table 2: Total magnetic moment ( $\mu_T$ ) of Nb-(HS)G/MoS<sub>2</sub> and Tc-(HS)G/MoS<sub>2</sub> materials are obtained by asymmetrically distributed up-spin and down-spin of electrons in 2s & 2p orbitals of C atoms; 4p, 4d & 5s orbitals of Mo atoms; 3s & 3p orbitals of S atoms; 4p, 4d & 5s orbitals of Nb atoms, and 4p, 4d & 5s orbitals of Tc atoms in the materials.**

Magnetic moments ( $\mu$ ) of (HS)G/MoS <sub>2</sub> , Nb-(HS)G/MoS <sub>2</sub> and Tc-(HS)G/MoS <sub>2</sub> materials	Nb-(HS)G/MoS <sub>2</sub> ( $\mu_B$ /cell)	Tc-(HS)G/MoS <sub>2</sub> ( $\mu_B$ /cell)
$\mu$ of C atoms by distributed spins in 2s orbital	0.00	0.00
$\mu$ of C atoms by distributed spins in 2p orbital	-0.01	0.01
$\mu$ of Mo atoms by distributed spins in 4p orbital	-0.03	0.00
$\mu$ of Mo atoms by distributed spins in 4d orbital	-0.10	0.02
$\mu$ of Mo atoms by distributed spins in 5s orbital	0.00	0.00
$\mu$ of S atoms by distributed spins in 3s orbital	0.00	0.00
$\mu$ of S atoms by distributed spins in 3p orbital	-0.08	0.04
$\mu$ of Nb atom by distributed spins in 4p orbital	-0.01	-
$\mu$ of Nb atom by distributed spins in 4d orbital	-0.01	-
$\mu$ of Nb atom by distributed spins in 5s orbital	0.00	-
$\mu$ of Tc atom by distributed spins in 4p orbital	-	0.00
$\mu$ of Tc atom by distributed spins in 4d orbital	-	0.01
$\mu$ of Tc atom by distributed spins in 5s orbital	-	0.00
Total magnetic moment ( $\mu_T$ ) $\mu_B$ /cell	-0.24	0.07

We observed that in DoS and PDoS plots of (HS)G/MoS<sub>2</sub> material, up-spin and down-spin states are symmetrically distributed near the Fermi energy level as shown in figs. 2(b-c). And we obtained the value of net magnetic moment given by up-spin down-spin of electrons in all individual orbital of C, Mo, & S atoms are of zero (0.00 $\mu_B$ /cell) value. We also calculated the magnetic moment due to integrated DoS and found that it is 0.00 $\mu_B$ /cell. Therefore, (HS)G/MoS<sub>2</sub> has non-magnetic properties.

Additionally, PDoS of up-spin and down-spin states of electrons near the Fermi level of Nb-(HS)G/MoS<sub>2</sub> are asymmetrically distributed as shown in figs. 3(b-c). The magnetic moment developed due to the distributed up-spin and down-spin in 2s & 2p orbitals of C atoms are 0.00  $\mu_B$ /cell & -0.01 $\mu_B$ /cell; 4p, 4d & 5s orbitals of Mo atoms are -0.03 $\mu_B$ /cell, -0.10  $\mu_B$ /cell & -0.00 $\mu_B$ /cell; 3s & 3p orbitals of S atoms are 0.00 $\mu_B$ /cell & -0.08  $\mu_B$ /cell; 4p, 4d & 5s orbitals of Nb atoms are -0.01 $\mu_B$ /cell, -0.01  $\mu_B$ /cell & 0.00  $\mu_B$ /cell values respectively. We obtained the total magnetic

moment of Nb-(HS)G/MoS<sub>2</sub> material is -0.24  $\mu_B$ /cell. Also, we have estimated the magnetic moment of Nb-(HS)G/MoS<sub>2</sub> material based on integrated density of states (IDoS) is 0.26  $\mu_B$ /cell. Therefore, Nb-(HS)G/MoS<sub>2</sub> material is called magnetic material. PDoS of up-spin and down-spin states of electrons are asymmetrically distributed around the Fermi energy level of Tc-(HS)G/MoS<sub>2</sub> material as shown in figs. 4(b-c). Magnetic moment developed in material due to unpaired up-spin and down-spin of electrons in 2s & 2p orbitals of C atoms are 0.00  $\mu_B$ /cell & 0.01  $\mu_B$ /cell; 4p, 4d & 5s orbitals of Mo atoms are 0.00  $\mu_B$ /cell, 0.02  $\mu_B$ /cell & 0.00  $\mu_B$ /cell; 3s & 3p orbitals of S atoms are 0.00  $\mu_B$ /cell & 0.04  $\mu_B$ /cell; 4p, 4d & 5s orbitals of Tc atoms are 0.00  $\mu_B$ /cell, 0.01  $\mu_B$ /cell & 0.00  $\mu_B$ /cell values respectively. The calculated value of total magnetic moment of Tc-(HS)G/MoS<sub>2</sub> material is 0.07  $\mu_B$ /cell. We calculated the magnetic moment of Tc-(HS)G/MoS<sub>2</sub> based on IDoS, which is 0.08  $\mu_B$ /cell value. Hence, Tc-(HS)G/MoS<sub>2</sub> material is also called magnetic material. Therefore, non-magnetic (HS)G/MoS<sub>2</sub> changes to magnetic Nb-

(HS)G/MoS<sub>2</sub> and Tc-(HS)G/MoS<sub>2</sub> materials due to Nb and Tc impurity defects respectively. In impurity defected Nb-(HS)G/MoS<sub>2</sub> material, major contribution of magnetic moment is given by distributed spin states in 4p, 4d orbitals of Mo atoms and 3p orbital of S atoms, while distributed spin states in 4d orbital of Mo atoms and 3p orbital of S atoms have dominant role for the development of magnetic moment in Tc-(HS)G/MoS<sub>2</sub> material.

#### 4. CONCLUSIONS

Structural, electronic and magnetic properties of (HS)G/MoS<sub>2</sub>, Nb-(HS)G/MoS<sub>2</sub> and Tc-(HS)G/MoS<sub>2</sub> materials have been studied using spin-polarized DFT with vdW corrections DFT-D2 approach. By analyzing the structures, we found that (HS)G/MoS<sub>2</sub> is more compact than impurity defected Nb-(HS)G/MoS<sub>2</sub> and Tc-(HS)G/MoS<sub>2</sub> materials. From the band structure calculations, we found that all the materials have metallic properties. It is found that n-type Schottky contact is formed in pristine (HS)G/MoS<sub>2</sub> and Tc impurity defect Tc-(HS)G/MoS<sub>2</sub> materials, and p-type Schottky contact is formed in Nb impurity defect Nb-(HS)G/MoS<sub>2</sub> material. Therefore, n-type Schottky contact of (HS)G/MoS<sub>2</sub> changes to p-type Schottky contact due to Nb impurity defect in (HS)G/MoS<sub>2</sub> material. To investigate the magnetic properties, we have carried out DoS and PDoS calculations, and found that (HS)G/MoS<sub>2</sub> is non-magnetic material, while Nb-(HS)G/MoS<sub>2</sub> and Tc-(HS)G/MoS<sub>2</sub> are magnetic materials. The total magnetic moment of Nb-(HS)G/MoS<sub>2</sub> and Tc-(HS)G/MoS<sub>2</sub> have values - 0.24 $\mu_B$ /cell and 0.07  $\mu_B$ /cell respectively. High value of magnetic moment is given by distributed unpaired up-spin and down-spin of electrons in 4p & 4d orbitals of Mo atoms; 3p orbital of S atoms in Nb-(HS)G/MoS<sub>2</sub> material, and 4d orbital of Mo atoms; 3p orbital of S atoms in Tc-(HS)G/MoS<sub>2</sub> material.

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